

# فیزیک 1

درس 24

صحرائی

گروه فیزیک دانشگاه رازی

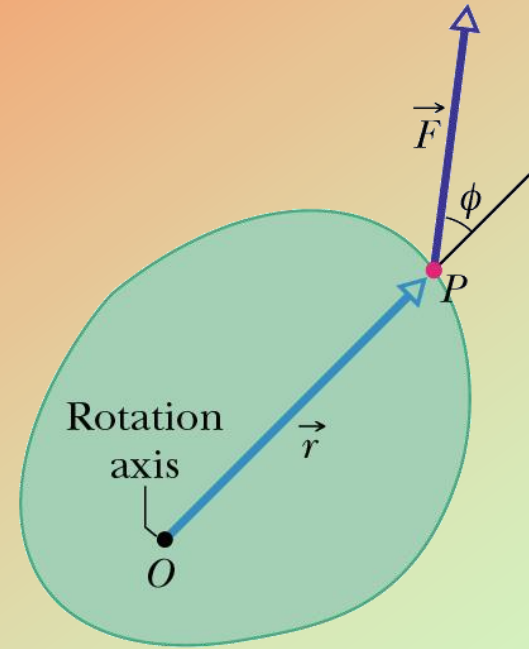
<http://www.razi.ac.ir/sahraei>

# Torque

- **When you try to swing a door, it's clear that it is easier to move the door if you apply the force farther away from the axis of rotation (e.g., the hinges).**
- **It's also clear that the application of the force is most efficient if the force is applied perpendicular to the plane of the door.**

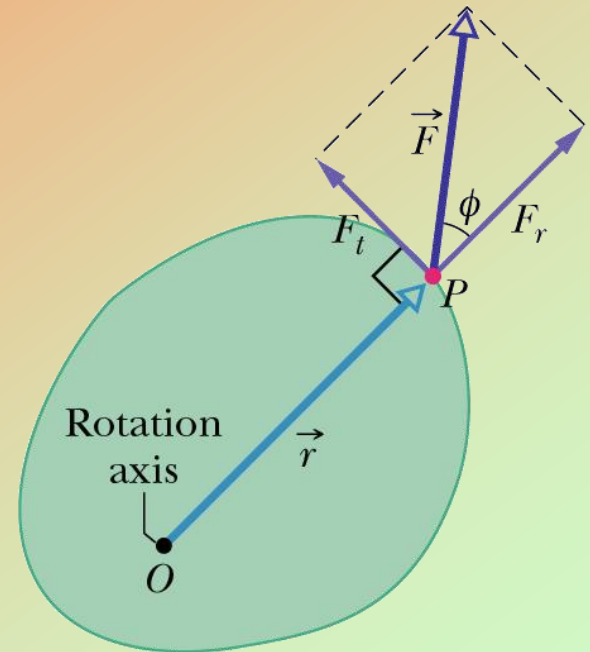
# Torque

- Lets look at an object that is rotating about a fixed axis  $O$
- A force  $F$  is applied at point  $P$  which is at a position  $r$  relative to the axis  $O$
- Note also that the force  $F$  is applied at an angle  $\Phi$  relative to the vector  $r$
- For simplicity we also assume that the force  $F$  is in the plane of the screen



# Torque

- Lets decompose the force into its components relative to the vector  $r$  :
  - The radial component is labeled  $F_r$  ( $F\cos\Phi$ )
  - The tangential component is labeled  $F_t$  ( $F\sin\Phi$ )



# Torque

- We therefore define *torque* to be the product of these two values:

$$\tau = (r)(F \sin \phi)$$

- Torque comes from the Latin word meaning “to twist”

## Torque

- By rearranging things a bit in the previous equation we can equivalently see that:

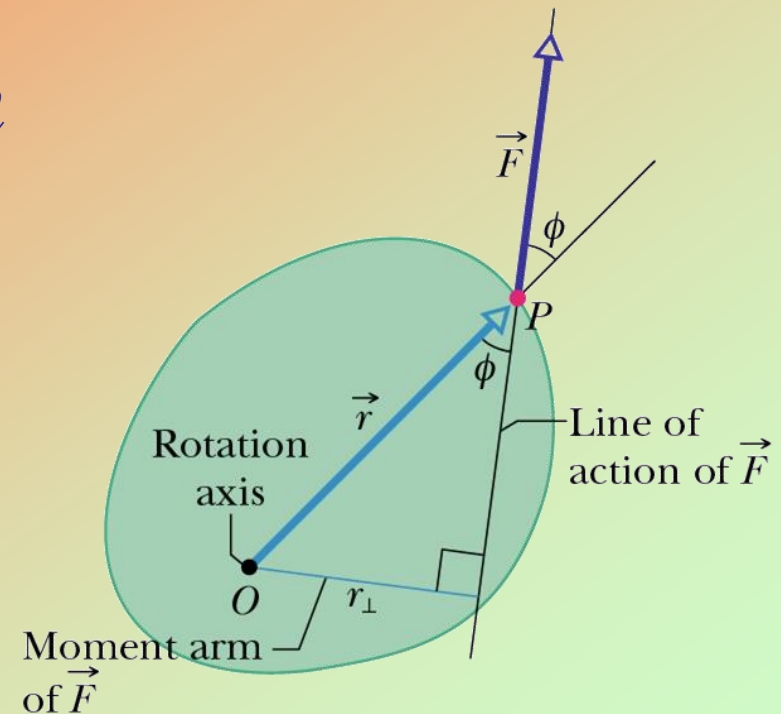
$$\tau = (r)(F \sin \phi) = rF_t$$

$$\tau = (r \sin \phi)(F) = r_{\perp} F$$

where  $r_{\perp}$  is the perpendicular distance from the rotation axis  $O$  to an extended line running through the vector  $F$  at point  $P$

# Torque

- The extended line is called the *line of action* and the value  $r_{\perp}$  is called the *moment arm*.
- Clearly if the force is applied completely tangentially, then the moment arm is just  $r$ .



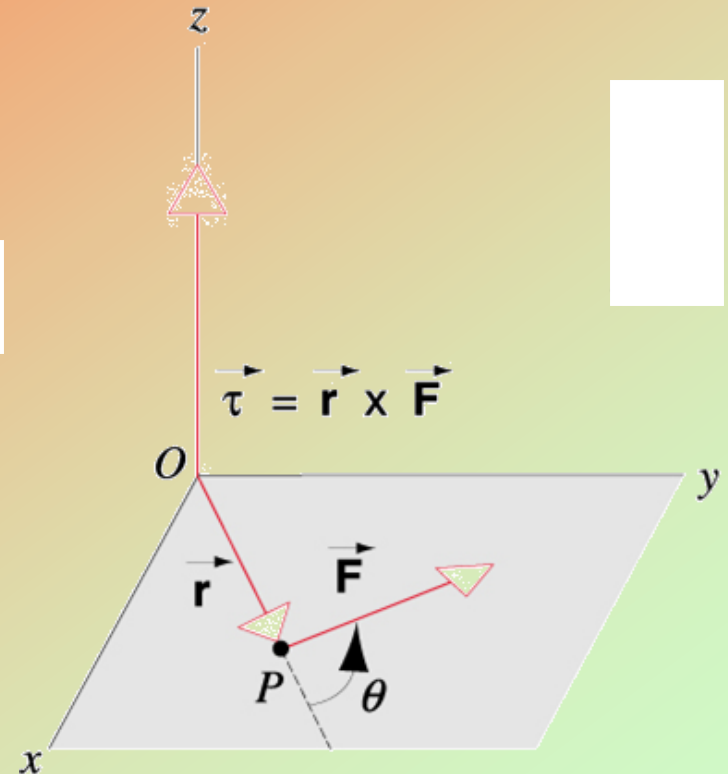
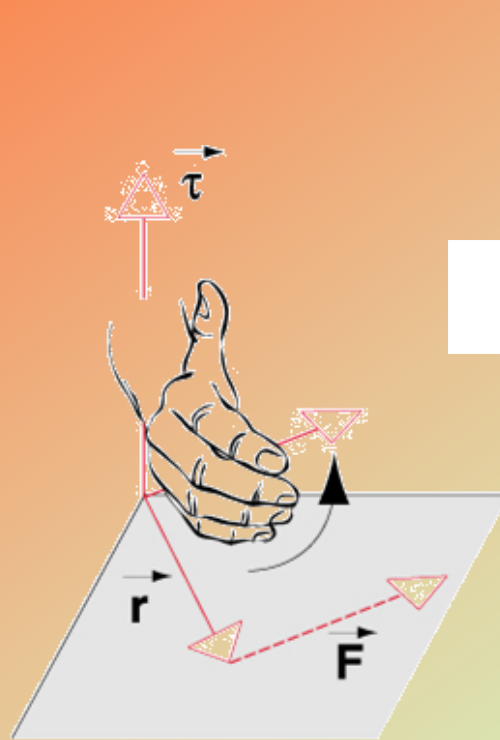
If an object rotates counterclockwise then the torque is positive – and vice versa (remember that clocks are negative...)

$$\vec{\tau} = \vec{r} \times \vec{F}$$

$$\tau = rF\sin\theta$$

*N.m , lb.ft*

$$ML^2T^{-2}$$



09.04

Work can also be expressed in joules ( $1 \text{ J} = 1 \text{ N}\cdot\text{m}$ ), but torque is never expressed that way.

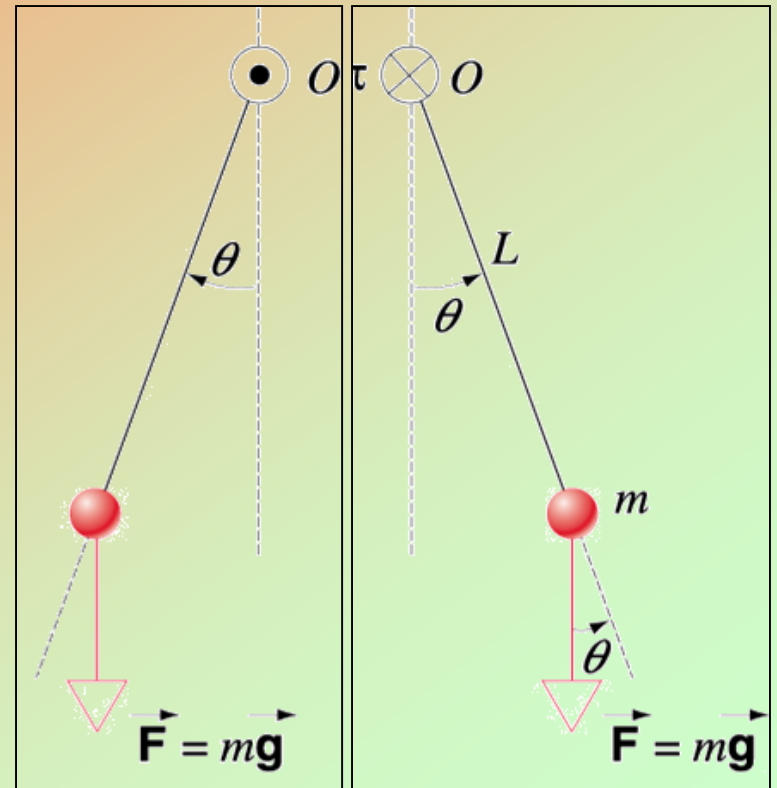
مثال: آونگی تشکیل شده است از جسمى به جرم  $m=0.17\text{ kg}$  که به انتهای میله صلبی به طول  $L=1.25\text{ m}$  و جرم ناچیز متصل است. مقدار گشتاور ناشی از گرانش حول نقطه  $O$  در لحظه ای که آونگ به اندازه  $\theta=10^\circ$  از امتداد قائم منحرف شده چقدر است؟

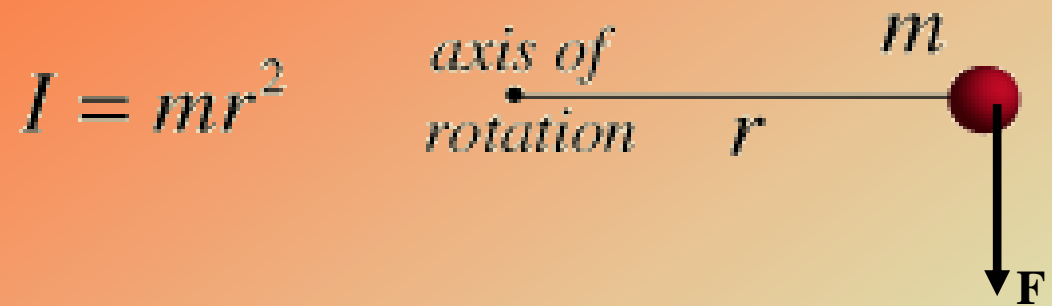
(ب) جهت گشتاور حول نقطه  $O$  در آن لحظه کدام است؟ آیا جهت گشتاور بستگی دارد به اینکه آونگ به کدام طرف خط قائم، چپ یا راست، جابجا می شود؟

$$\vec{\tau} = \vec{r} \times \vec{F} = rF \sin \theta$$

$$\tau = Lmg \sin \theta$$

$$= (1.25\text{m})(0.17\text{kg})(9.8\text{m/s}^2)(\sin 10^\circ) = 0.36\text{N.m}$$





$$\vec{\tau} = \vec{r} \times \vec{F} = rF \sin \theta$$

$$\tau = mar \qquad \tau = I\alpha$$

$$mar = I\alpha \qquad a_T = r\alpha$$

$$I = mr^2$$

# Work and Rotational Kinetic Energy

Suppose that the change in the kinetic energy is the only change in the overall energy of the system – thus:

$$\Delta K = K_f - K_i = \frac{1}{2}mv_f^2 - \frac{1}{2}mv_i^2 = W$$

For motion confined to a single axis (let's say the  $x$  axis) we have:

$$W = \int_{x_i}^{x_f} F dx \qquad W = F(x_f - x_i)$$

And finally, the power is:  $P = \frac{dW}{dt} = Fv$

In this case suppose that the change in the rotational kinetic energy is the only change in the overall energy of the system – thus:

$$\Delta K = K_f - K_i = \frac{1}{2} I \omega_f^2 - \frac{1}{2} I \omega_i^2 = W$$

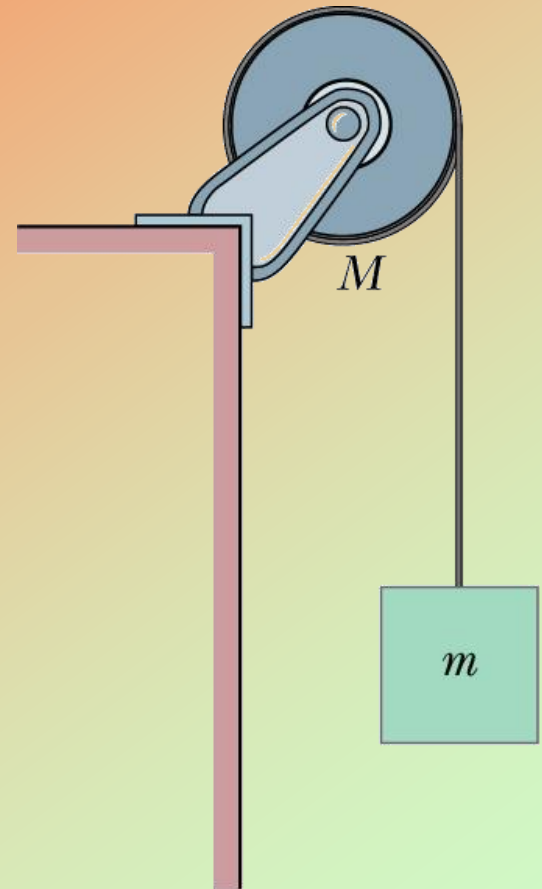
The corresponding equation for work in the rotational case is:

$$W = \int_{\theta_i}^{\theta_f} \tau d\theta \qquad W = \tau(\theta_f - \theta_i)$$

when the torque is constant and the angular displacement goes from  $\theta_i$  to  $\theta_f$

$$P = \frac{dW}{dt} = \tau\omega$$

- We have a disk with mass  $M = 2.5 \text{ kg}$  and a radius  $R = 20 \text{ cm}$  mounted on a fixed, frictionless axle
- A block of mass  $m = 1.2 \text{ kg}$  hangs from a massless string which is wrapped around the disk several times
- Find the acceleration of the block, the angular acceleration of the disk and the string tension



- As usual, we begin by drawing the free body diagrams
- Starting with the block we see that our force equation turns out to be:

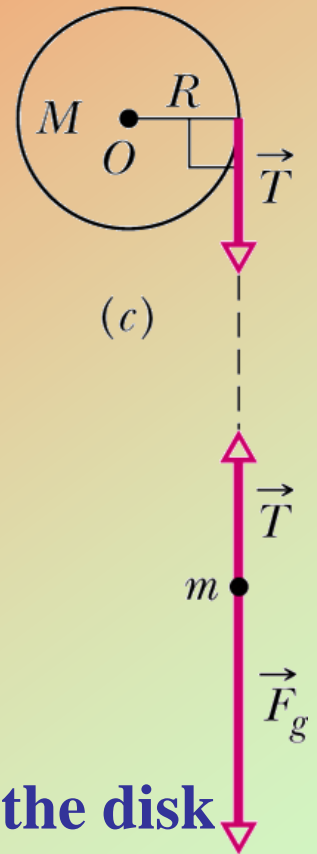
$$T - mg = ma$$

- We have the acceleration  $a$  in this equation, but we can't solve for it yet as we don't know the value of  $T$

- Moving on to the disk, we can see that the torque on the disk is:

$$\tau = -RT$$

- (since the disk is turning clockwise, the torque is negative)



$$\tau_{net} = I\alpha$$

$$I = \frac{1}{2}MR^2$$

$$-RT = \frac{1}{2}MR^2\alpha$$

$$a_T = \alpha R$$

$$\alpha = \frac{a}{R}$$

$$T = -\frac{1}{2}Ma$$

$$T - mg = ma$$

$$a = -g \frac{2m}{M + 2m} = -4.8 \text{ m/s}^2$$

$$T = -\frac{1}{2}Ma = 6.0 \text{ N}$$

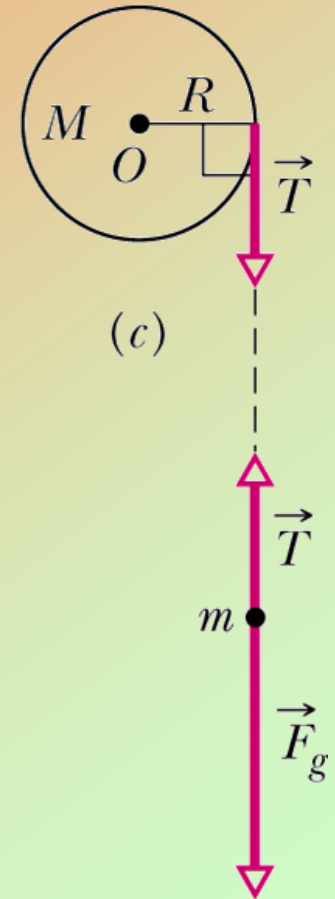
$$\alpha = \frac{a}{R} = \frac{-4.8 \text{ m/s}^2}{0.2 \text{ m}} = -24 \text{ rad/s}^2 = -3.8 \text{ rev/s}^2$$

As a final check let's see what happens when  $M = 0$ :

$$a = -g \frac{2m}{M + 2m} = -g$$

$$T = \frac{1}{2}Ma = 0$$

مثال فوق را از دیدگاه کار و انرژی بررسی کنید.



$$W_{net} = mgL$$

$$\Delta K = K_f - K_i = K_f = \frac{1}{2} I \omega^2 + \frac{1}{2} m v^2$$

$$W_{net} = \Delta K$$

$$mgL = \frac{1}{2} I \omega^2 + \frac{1}{2} m v^2 = \frac{1}{2} \left( \frac{1}{2} M R^2 \right) \left( \frac{v}{R} \right)^2 + \frac{1}{2} m v^2$$

$$v^2 = 2 \left[ \frac{2m g}{M + 2m} \right] L \qquad v^2 = v_0^2 + 2ax$$

$$W_{net} = mgL - TL$$

$$mgL - TL = \frac{1}{2} m v^2$$

$$W_{net} = TR\phi = TL$$

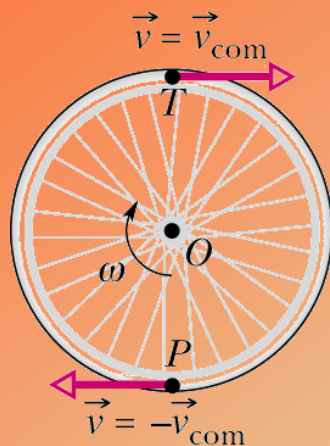
$$\Delta K = \frac{1}{2}I\omega^2 = \frac{1}{2}\left(\frac{1}{2}MR^2\right)\frac{v^2}{R^2} = \frac{1}{4}Mv^2$$

$$TL = \frac{1}{4}Mv^2$$

# Rolling

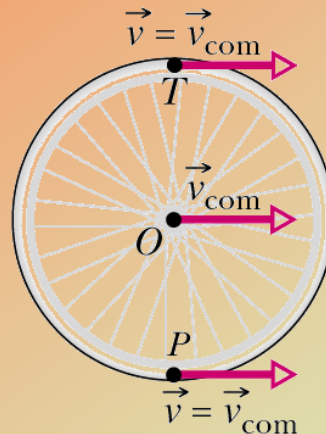
Here we can see that rolling motion is a combination of purely rotational and purely translational motions

(a) Pure rotation



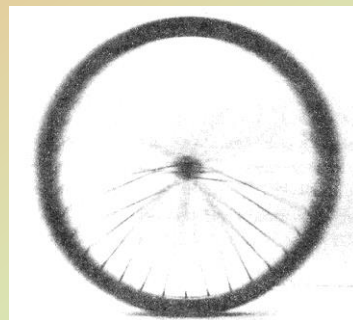
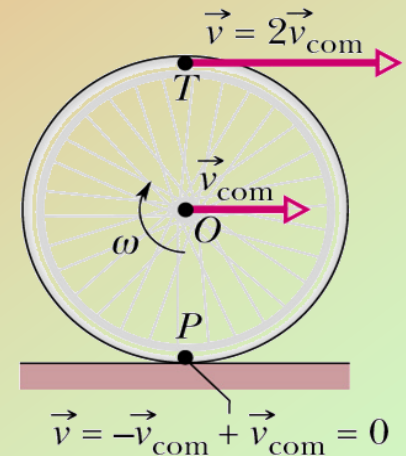
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(b) Pure translation

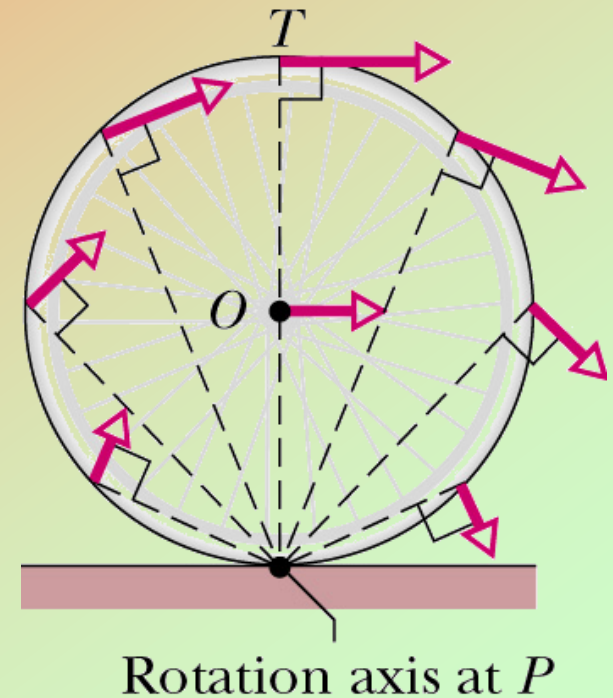


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(c) Rolling motion



- There is another way to look at this however – a way which sees rolling as purely rotational
- In this case we view the rotation as being around the point where the wheel contacts the ground...
- The rotation axis is taken as the point  $P$
- The vectors in the figure represent the instantaneous velocity of various points on the rolling wheel

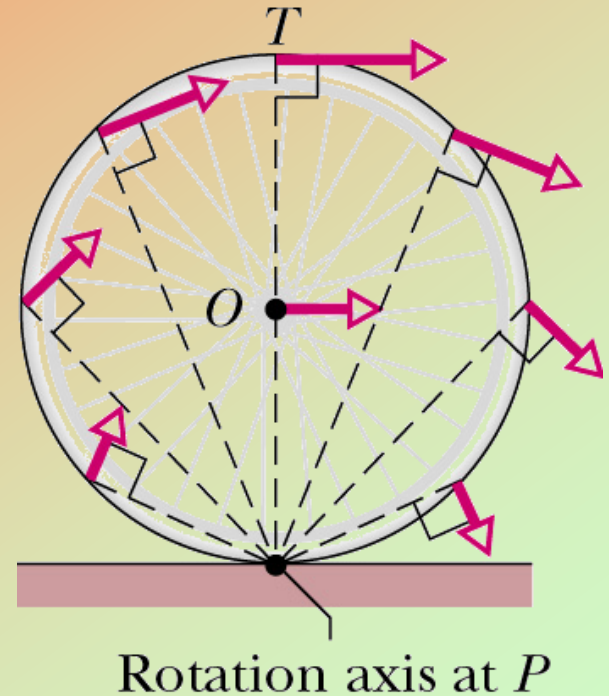


# The Kinetic Energy of Rolling

- We want to calculate the kinetic energy of a wheel rotating about the axis  $P$
- We know that (since this is a purely rotational view of the problem) we have:

$$K = \frac{1}{2} I_P \omega^2 \quad I_P = I_{cm} + MR^2$$

$$K = \frac{1}{2} (I_{cm} + MR^2) \omega^2$$



# The Kinetic Energy of Rolling

- Substituting for  $I_P$  in our equation for  $K$  and then expanding it out we get:

$$K = \frac{1}{2} I_{cm} \omega^2 + \frac{1}{2} MR^2 \omega^2$$

- and when we also substitute in  $v_{cm} = \omega R$  we get:

$$K = \frac{1}{2} I_{cm} \omega^2 + \frac{1}{2} M v_{cm}^2$$

## The Kinetic Energy of Rolling

**But this is exactly equivalent to the kinetic energy of a wheel rotating about its CM added to the kinetic energy of a body's CM in translation:**

$$K = \frac{1}{2} I_{cm} \omega^2 + \frac{1}{2} M v_{cm}^2$$

**Rotational  
Kinetic Energy**

**Translational  
Kinetic Energy**

**A uniform solid cylindrical disk, of mass  $M = 1.4$  kg and radius  $R = 8.5$  cm, rolls smoothly across a horizontal table at a speed of 15 cm/s.**

**What is the kinetic energy  $K$  of the disk?**

$$K = \frac{1}{2} I_{cm} \omega^2 + \frac{1}{2} M v_{cm}^2$$

**speed of the center of mass So  $v_{cm} = 15$  cm/s = 0.15 m/s**

$$\omega = \frac{v_{cm}}{R} \quad I_{cm} = \frac{1}{2} MR^2$$

$$K = \frac{1}{2} I_{cm} \omega^2 + \frac{1}{2} M v_{cm}^2$$

$$K = \frac{1}{2} \left( \frac{1}{2} M R^2 \right) \left( \frac{v_{cm}}{R} \right)^2 + \frac{1}{2} M v_{cm}^2$$

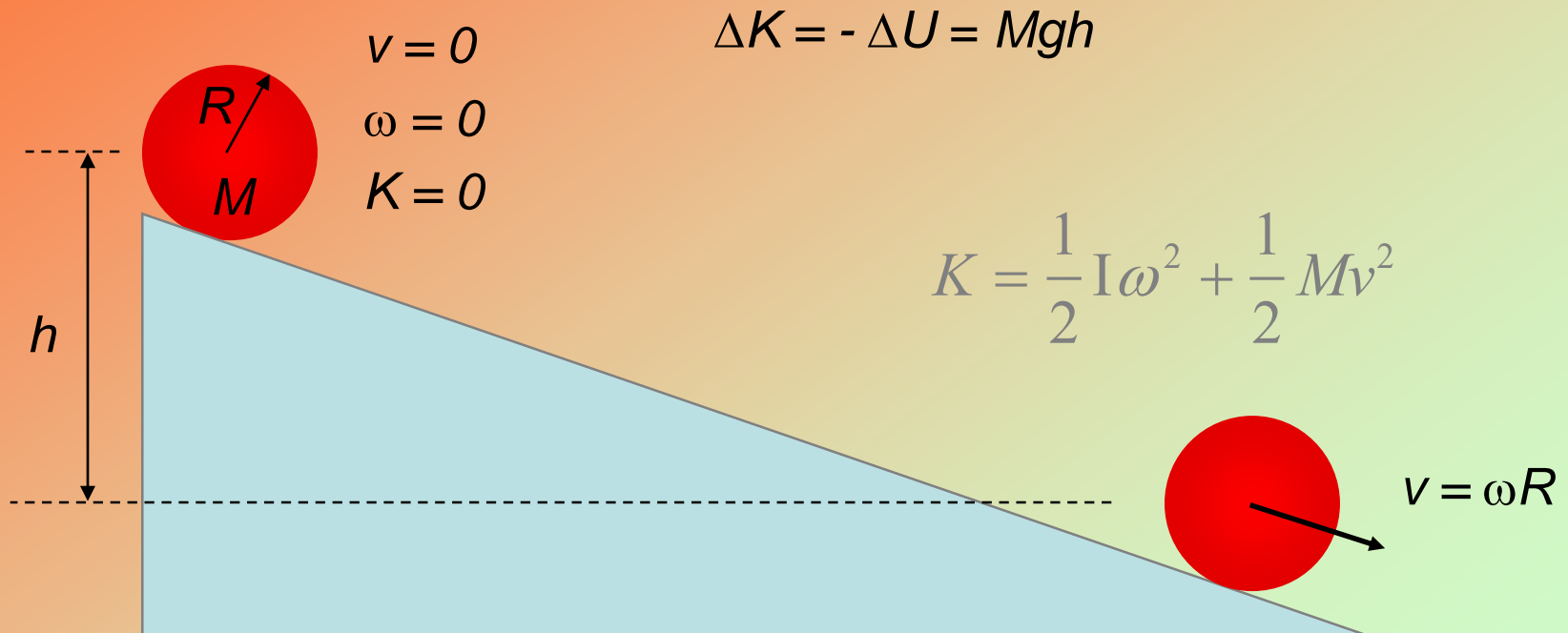
$$K = \frac{3}{4} M v_{cm}^2$$

It doesn't depend on the radius of the disk

$$K = \frac{3}{4} (1.4 \text{ kg}) (0.15 \text{ m/s})^2 = 0.024 \text{ J}$$

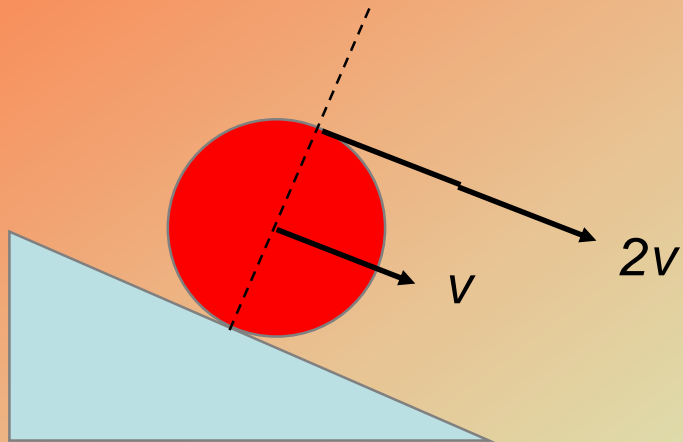
# Rolling Motion

**Objects of different  $I$  rolling down an inclined plane:**

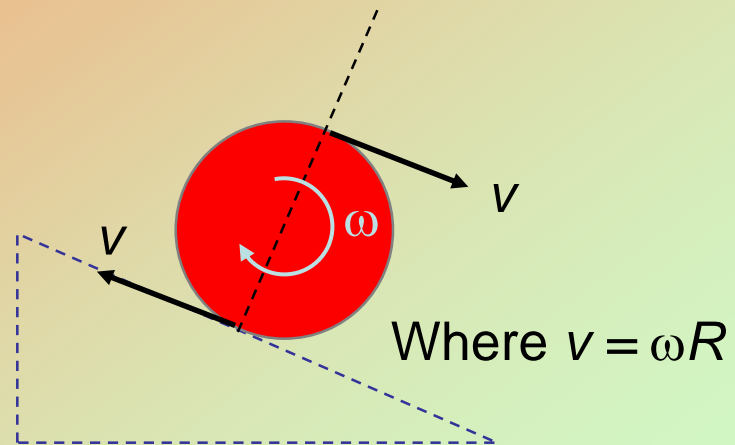


# Rolling...

- If there is no slipping:



In the lab reference frame



In the CM reference frame

# Rolling...

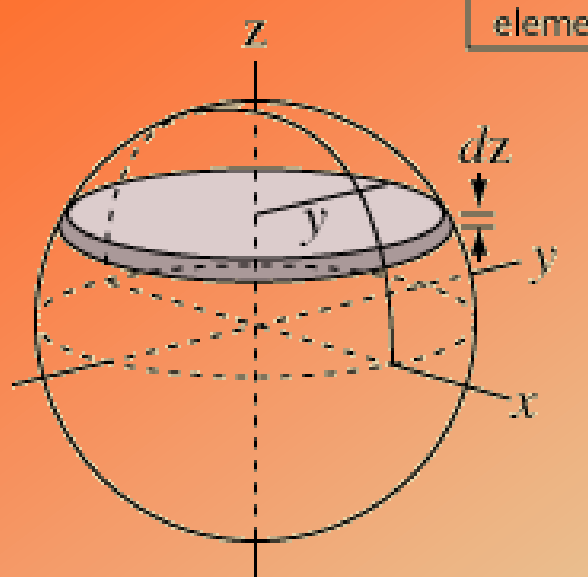
$$K = \frac{1}{2}I\omega^2 + \frac{1}{2}Mv^2 \quad \text{Use } v = \omega R \text{ and } I = cMR^2.$$

hoop:  $c = 1$   
disk:  $c = 1/2$   
sphere:  $c = 2/5$   
etc...

So:  $\frac{1}{2}(c+1)Mv^2 = Mgh \quad \Rightarrow \quad v = \sqrt{2gh} \sqrt{\frac{1}{c+1}}$

The rolling speed is always lower than in the case of simple sliding since the kinetic energy is shared between CM motion and rotation.

# Moment of Inertia: Sphere



Infinitesimal moment of inertia element

Moment of a thin disk of mass  $dm$

Substituting area times height of disk for volume  $dV$

$$dI = \frac{1}{2} y^2 dm = \frac{1}{2} y^2 \rho dV = \frac{1}{2} y^2 \rho \pi y^2 dz$$

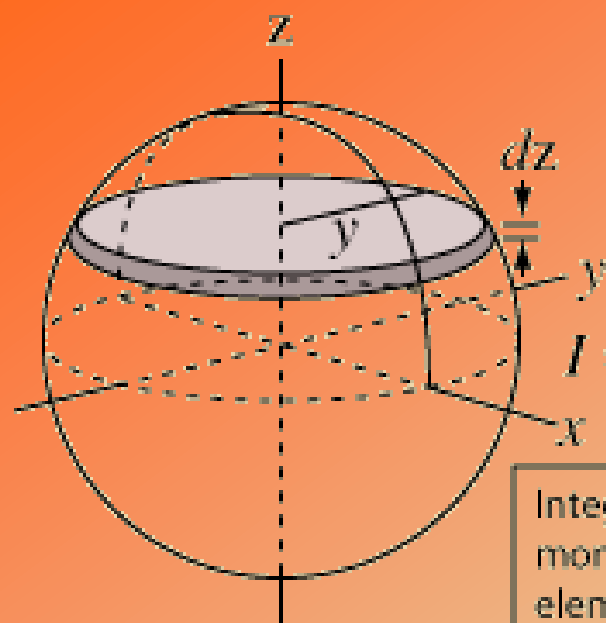
Integrating the moment of inertia element from the bottom to the top of the sphere

Expressing mass  $dm$  in terms of density and volume.

$$I = \frac{1}{2} \rho \pi \int_{-R}^R y^4 dz = \frac{1}{2} \rho \pi \int_{-R}^R (R^2 - z^2)^2 dz = \frac{8}{15} \rho \pi R^5$$

Radius =  $R$   
 Mass =  $M$   
 Density =  $\rho = \frac{M}{V} = \frac{M}{\frac{4}{3} \pi R^3}$

$$I = \frac{8}{15} \left[ \frac{M}{\frac{4}{3} \pi R^3} \right] \pi R^5 = \frac{2}{5} MR^2$$



Find  $y$  in terms of the variable of integration  $z$



$$I = \frac{1}{2} \rho \pi \int_{-R}^R y^4 dz = \frac{1}{2} \rho \pi \int_{-R}^R (R^2 - z^2)^2 dz = \frac{8}{15} \rho \pi R^5$$

Integrating the moment of inertia element from the bottom to the top of the sphere

$$(R^2 - z^2)^2 = R^4 - 2R^2z^2 + z^4$$

Polynomial form integral

Radius =  $R$

Mass =  $M$

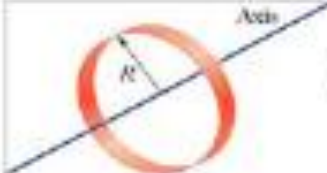


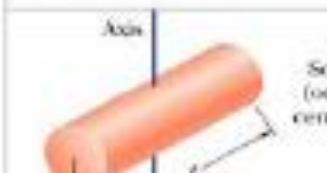
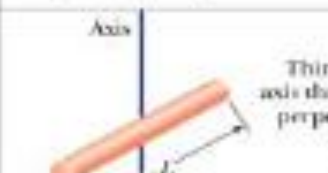
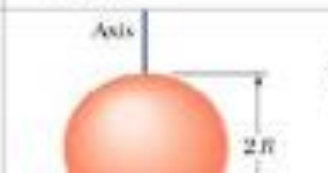

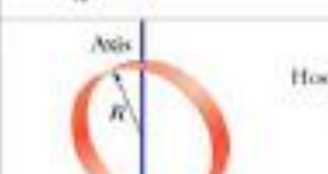
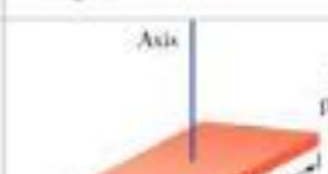
Density

$$= \rho = \frac{M}{V} = \frac{M}{\frac{4}{3} \pi R^3}$$

$$\int_{-R}^R \dots dz = \left[ R^4 z - \frac{2R^2 z^3}{3} + \frac{z^5}{5} \right]_{-R}^R = 2R^5 \left( \frac{15}{15} - \frac{10}{15} + \frac{3}{15} \right)$$

$$I = \frac{8}{15} \rho \pi R^5$$

- Rotational inertia involves not only the mass but also the distribution of mass for continuous masses
- Calculating the rotational inertia  $I = \int r^2 dm$

 <p>Hoop about central axis</p> <p><math>I = MR^2</math> (a)</p>	 <p>Annular cylinder (or ring) about central axis</p> <p><math>I = \frac{1}{2} M(R_1^2 + R_2^2)</math> (b)</p>	 <p>Solid cylinder (or disk) about central axis</p> <p><math>I = \frac{1}{2} MR^2</math> (c)</p>
 <p>Solid cylinder (or disk) about central diameter</p> <p><math>I = \frac{1}{2} MR^2 + \frac{1}{12} ML^2</math> (d)</p>	 <p>Thin rod about axis through center perpendicular to length</p> <p><math>I = \frac{1}{12} ML^2</math> (e)</p>	 <p>Solid sphere about any diameter</p> <p><math>I = \frac{2}{5} MR^2</math> (f)</p>
 <p>Thin spherical shell about any diameter</p> <p><math>I = \frac{2}{3} MR^2</math> (g)</p>	 <p>Hoop about any diameter</p> <p><math>I = \frac{1}{2} MR^2</math> (h)</p>	 <p>Slab about perpendicular axis through center</p> <p><math>I = \frac{1}{12} M(a^2 + b^2)</math> (i)</p>