# Dynamic Meteorology 1

Lecture 7

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#### Total Differentiation of a Vector in a Rotating Frame of Reference

Before we can write Newton's second law of motion for a reference frame rotating with the earth, we need to develop a relationship between the total derivative of a vector in an inertial reference frame and the corresponding derivative in a rotating system.

Let A be an arbitrary vector with Cartesian components

 $\vec{A} = A_x \hat{i} + A_y \hat{j} + A_z \hat{k}$  in an inertial frame of reference, and

 $\vec{A} = A'_x \hat{i}' + A'_y \hat{j}' + A'_z \hat{k}'$  in a rotating frame of reference

ماهیت بردار در دستگاه لخت و چرخان یکی است.

 $\vec{A} = \begin{vmatrix} x & x' \\ y & y' \\ z & z' \end{vmatrix}$ 

α

 $\vec{A} = A_x \hat{i} + A_y \hat{j} + A_z \hat{k}$ 

in an inertial frame of reference, then

 $\frac{d\vec{A}}{dt} = \left(A_{x}\frac{d\hat{i}}{dt} + \hat{i}\frac{dA_{x}}{dt}\right) + \left(A_{y}\frac{d\hat{j}}{dt} + \hat{j}\frac{dA_{y}}{dt}\right) + \left(A_{z}\frac{d\hat{k}}{dt} + \hat{k}\frac{dA_{z}}{dt}\right)$ 

Since the coordinate axes are in an inertial frame of reference,

$$\frac{d\hat{i}}{dt} = \frac{d\hat{j}}{dt} = \frac{d\hat{k}}{dt} = 0$$

$$\frac{dA}{dt} = \frac{dA_x}{dt}\hat{i} + \frac{dA_y}{dt}\hat{j} + \frac{dA_z}{dt}\hat{k} \qquad (Eq. 1)$$

If  $\vec{A} = A'_x \hat{i}' + A'_y \hat{j}' + A'_z \hat{k}'$  in a rotating frame of reference, then

 $\frac{d\vec{A}}{dt} = \left(A'_x \frac{d\hat{i}'}{dt} + \hat{i}' \frac{dA'_x}{dt}\right) + \left(A'_y \frac{d\hat{j}'}{dt} + \hat{j}' \frac{dA'_y}{dt}\right) + \left(A'_z \frac{d\hat{k}'}{dt} + \hat{k}' \frac{dA'_z}{dt}\right)$ (Eq. 2)

Because the left hand sides of (Eq. 1) and (Eq. 2) are identical,

 $\frac{dA_x}{dt}\hat{i} + \frac{dA_y}{dt}\hat{j} + \frac{dA_z}{dt}\hat{k} = \left(A_x'\frac{d\hat{i}'}{dt} + \hat{i}'\frac{dA_x'}{dt}\right) + \left(A_y'\frac{d\hat{j}'}{dt} + \hat{j}'\frac{dA_y'}{dt}\right) + \left(A_z'\frac{d\hat{k}'}{dt} + \hat{k}'\frac{dA_z'}{dt}\right)$ 

Regrouping the terms

 $\frac{dA_x}{dt}\hat{i} + \frac{dA_y}{dt}\hat{j} + \frac{dA_z}{dt}\hat{k} = \frac{dA'_x}{dt}\hat{i}' + \frac{dA'_y}{dt}\hat{j}' + \frac{dA'_z}{dt}\hat{k}' + A'_x\left(\frac{d\hat{i}'}{dt}\right) + A'_y\left(\frac{d\hat{j}'}{dt}\right) + A'_z\left(\frac{d\hat{k}'}{dt}\right)$  $\left(rac{dA}{dt}
ight)_{inertial}$  $\left(\frac{dA}{dt}\right)$ effects of rotation

To interpret

 $\frac{d\hat{i}'}{dt}, \quad \frac{d\hat{j}'}{dt}, \quad \frac{d\hat{k}'}{dt}$ 

#### think of each unit vector as a position vector

### linear velocity = angular velocity x position vector

 $\vec{V} = \vec{\Omega} \times \vec{r}$ 

Because

 $\vec{V} = \frac{d\vec{r}}{dt}, \qquad \frac{d\vec{r}}{dt} = \vec{\Omega} \times \vec{r}$ 

Thus

 $\frac{d\hat{i}'}{dt} = \vec{\Omega} \times \hat{i}', \qquad \frac{d\hat{j}'}{dt} = \vec{\Omega} \times \hat{j}', \qquad \frac{d\hat{k}'}{dt} = \vec{\Omega} \times \hat{k}'$ 





(effects of rotation)

 $\vec{\Omega} \times \vec{A}$ 

 $\frac{dA_x}{dt}\hat{i} + \frac{dA_y}{dt}\hat{j} + \frac{dA_z}{dt}\hat{k} = \frac{dA'_x}{dt}\hat{i}' + \frac{dA'_y}{dt}\hat{j}' + \frac{dA'_z}{dt}\hat{k}' + A'_x(\vec{\Omega}\times\hat{i}') + A'_y(\vec{\Omega}\times\hat{j}') + A'_z(\vec{\Omega}\times\hat{k}')$ 



This equation provides us with a formal way of expressing the balance of forces on a fluid parcel in a rotating coordinate system.

#### Newton's second law in an inertial reference frame:



To transform to rotating coordinates:



 $\vec{r}$  is the position vector for an air parcel on the rotating earth

 $\vec{V}_{inertial} = \vec{V} + \vec{\Omega} \times \vec{r}$ 

Velocity is the rate of change of the position vector with time



 $\left(\frac{d\vec{V}_{inertial}}{dt}\right) = \frac{d\vec{V}_{inertial}}{dt} + \vec{\Omega} \times \vec{V}_{inertial}$ 

Using the transformation of the total derivative

 $\left(\frac{dV_{inertial}}{dt}\right)_{inertial} = \frac{dV_{inertial}}{dt} + \vec{\Omega} \times \vec{V}_{inertial}$ 



 $\left(\frac{dV_{inertial}}{dt}\right)_{inertial} = \frac{d}{dt}\left(\vec{V} + \vec{\Omega} \times \vec{r}\right) + \vec{\Omega} \times \left(\vec{V} + \vec{\Omega} \times \vec{r}\right)$ 

Using some vector identities and  $\left(\frac{d\vec{V}_{inertial}}{dt}\right)_{inertial} = \frac{d\vec{V}}{dt} + 2\vec{\Omega} \times \vec{V} - \Omega^2 \vec{R}$   $\int_{inertial} \frac{d\vec{V}}{dt} + 2\vec{\Omega} \times \vec{V} - \Omega^2 \vec{R}$   $\int_{inertial} \frac{d\vec{V}}{dt} + 2\vec{\Omega} \times \vec{V} - \Omega^2 \vec{R}$   $\int_{inertial} \frac{d\vec{V}}{dt} + 2\vec{\Omega} \times \vec{V} - \Omega^2 \vec{R}$   $\int_{inertial} \frac{d\vec{V}}{dt} + 2\vec{\Omega} \times \vec{V} - \Omega^2 \vec{R}$   $\int_{inertial} \frac{d\vec{V}}{dt} + 2\vec{\Omega} \times \vec{V} - \Omega^2 \vec{R}$   $\int_{inertial} \frac{d\vec{V}}{dt} + 2\vec{\Omega} \times \vec{V} - \Omega^2 \vec{R}$   $\int_{inertial} \frac{d\vec{V}}{dt} + 2\vec{\Omega} \times \vec{V} - \Omega^2 \vec{R}$   $\int_{inertial} \frac{d\vec{V}}{dt} + 2\vec{\Omega} \times \vec{V} - \Omega^2 \vec{R}$   $\int_{inertial} \frac{d\vec{V}}{dt} + 2\vec{\Omega} \times \vec{V} - \Omega^2 \vec{R}$   $\int_{inertial} \frac{d\vec{V}}{dt} + 2\vec{\Omega} \times \vec{V} - \Omega^2 \vec{R}$   $\int_{inertial} \frac{d\vec{V}}{dt} + 2\vec{\Omega} \times \vec{V} - \Omega^2 \vec{R}$   $\int_{inertial} \frac{d\vec{V}}{dt} + 2\vec{\Omega} \times \vec{V} - \Omega^2 \vec{R}$   $\int_{inertial} \frac{d\vec{V}}{dt} + 2\vec{\Omega} \times \vec{V} - \Omega^2 \vec{R}$ 

Acceleration following the motion in an inertial system

Coriolis acceleration

Rate of change of relative velocity following the relative motion in a rotating reference frame.

Centrifugal acceleration Substituting into Newton's second law:





If the real forces acting on a fluid parcel are the pressure gradient force, gravitation and friction, then

# $\frac{dV}{dt} = -2\vec{\Omega} \times \vec{V} - \frac{1}{\rho}\nabla p + \vec{g} + \vec{F}_r$

Rate of change of relative velocity following the relative motion in a rotating reference frame.

Coriolis acceleration f

Pressure gradient force (per unit mass)

Gravity term (gravitation + centrifugal)

Friction

Vector momentum equation in rotating coordinates

#### Momentum Equations in Spherical Coordinates

For a variety of reasons, it is useful to express the vector momentum equation for a rotating earth as a set of scalar component equations.

The use of latitude-longitude coordinates to describe positions on earth's surface makes it convenient to write the momentum equations in spherical coordinates.

The coordinate axes are  $(\lambda, \phi, z)$ 

Where,

 $\lambda$  is longitude,

 $\phi$  is latitude,



z is height.

#### Orientation of Coordinate Axes

The x- and y-axes are customarily defined to point east and north, respectively, such that

 $dx = Rd\lambda = a\cos\phi\,d\lambda$ 

and

 $dy = a d\phi$ 

Thus the horizontal velocity components are

 $u = \frac{dx}{dt},$ 

 $v = \frac{dy}{dt}$ 

The unit vectors in the spherical coordinate system are functions of position

$$\frac{d\vec{V}}{dt} = \hat{i}\frac{du}{dt} + \hat{j}\frac{dv}{dt} + \hat{k}\frac{dw}{dt} + u\frac{d\hat{i}}{dt} + v\frac{d\hat{j}}{dt} + w\frac{d\hat{k}}{dt}$$

A Complication of Spherical Coordinates

When the x and y coordinates are defined in this way, the coordinate system is not strictly Cartesian, because the directions of the unit vectors depend on their position on the earth's surface.

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