معادلات ديفر انسيل

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Review of Lesson 2

Nonhomogeneous Differential Equation

General Remarks on Solutions

 $F(x, y, y', y'', ..., y^{(n)}) = 0 \rightarrow y = y(x, c_1, c_2, ..., c_n)$

General Solution

Particular Solution

The geometric meaning of a solution

Families of Curves. Orthogonal Trajectories

Linear Equation Applications

Falling Bodies and Other Rate problems

F = ma, $mg = m\frac{d^2 y}{dt^2}$ $\frac{d^2 y}{dt^2} = g \qquad , \qquad \frac{dv}{dt} = g$



dv = gdt, $v = gt + c_1$ t=0, $v=v_0 \rightarrow c_1=v_0$ $v=gt+v_0$ $\frac{dy}{dt} = gt + v_0 \qquad y_g = \frac{1}{2}gt^2 + v_0t + c_2$

$$t = 0 , y = y_0 , c_2 = y_0$$

$$y_g = \frac{1}{2}gt^2 + v_0t + y_0 \quad v = gt + v_0$$

If the body falls from rest starting at y = 0
so that $v_0 = y_0 = 0$

$$y_p = \frac{1}{2}gt^2 \quad \text{and} \quad v_p = gt$$

$$v = \sqrt{2gy}$$

C

Retarded fall 100 kg Parachuter 150 kg Parachuter F = ma, $m\frac{d^2y}{dt^2} = mg - kv$ $\frac{d^2 y}{dt^2} = g - cv \quad , \ k/m = c$ F_{gra⊽} dv $\frac{dv}{dt} = g - cv$ $\frac{dv}{g - cv} = dt$ $-\frac{1}{2}\ln(g - cv) = t + c_1 \quad \ln(g - cv) = -ct - cc_1$ C

 $\ln(g-cv) = -ct + c_2$ $c_2 = -cc_1$ $t = 0 \leftrightarrow v = 0 \rightarrow c_2 = \ln g$ $\ln \frac{g - cv}{dt} = -ct \qquad g - cv = ge^{-ct}$ $v = \frac{g}{c}(1 - e^{-ct}) = \frac{mg}{k}(1 - e^{-\frac{\kappa}{m}t})$ if $t \to \infty$, $e^{-ct} = 0 \implies v_t \to \frac{mg}{l_r}$

 $\frac{dy}{dt} = \frac{mg}{k}(1 - e^{-\frac{k}{m}t}) \quad dy = \frac{mg}{k}dt - \frac{mg}{k}e^{-\frac{k}{m}t}dt$ $y = \frac{mg}{k}t + \frac{m^2g}{k^2}e^{-\frac{k}{m}t} + c_3$ $t = 0 \leftrightarrow y = 0 \rightarrow c_3 = -\frac{m^2 g}{k^2}$ $y = \frac{mg}{k}t + \frac{m^2g}{k^2}e^{-\frac{k}{m}t} - \frac{m^2g}{k^2}$ $y = \frac{mg}{k}t + \frac{m^2g}{k^2}(e^{-\frac{k}{m}t} - 1)$ y as a function of t



$$2ga(\cos\theta - \cos\alpha) = a^{2}(d\theta/dt)^{2}$$
$$dt = -\sqrt{\frac{a}{2g}} \frac{d\theta}{\sqrt{\cos\theta - \cos\alpha}}$$
$$T/4 \int_{0}^{T/4} dt = \int_{0}^{\alpha} \sqrt{\frac{a}{2g}} \frac{d\theta}{\sqrt{\cos\theta - \cos\alpha}}$$

$$\frac{T}{4} = \sqrt{\frac{a}{2g}} \int_{0}^{\alpha} \frac{d\theta}{\sqrt{\cos\theta - \cos\alpha}}$$

$$\begin{cases} \cos\theta = 1 - 2\sin^2\frac{\theta}{2} \\ \cos\alpha = 1 - 2\sin^2\frac{\alpha}{2} \\ T = 2\sqrt{\frac{a}{g}} \int_{0}^{\alpha} \frac{d\theta}{\sqrt{\sin^2(\alpha/2) - \sin^2(\theta/2)}} \\ \sin(\alpha/2) = k &, \quad \sin(\theta/2) = k\sin\varphi \\ \cos\frac{\theta}{2} = \sqrt{1 - \sin^2\frac{\theta}{2}} & \cos\frac{\theta}{2} = \sqrt{1 - k^2\sin^2\varphi} \\ \frac{1}{2}\cos\frac{\theta}{2}d\theta = k\cos\varphi d\varphi \rightarrow d\theta \end{cases}$$

C

$$d\theta = \frac{2k\cos\varphi d\varphi}{\cos(\theta/2)} = \frac{2k\cos\varphi d\varphi}{\sqrt{1-k^2\sin^2\varphi}}$$
$$T = 2\sqrt{\frac{a}{g}} \int_0^{\alpha} \frac{d\theta}{\sqrt{\sin^2(\alpha/2) - \sin^2(\theta/2)}}$$
$$T = 4\sqrt{\frac{a}{g}} \int_0^{\pi/2} \frac{k\cos\varphi d\varphi}{\sqrt{1-k^2\sin^2\varphi}} \times \frac{1}{\sqrt{k^2 - k^2\sin^2\varphi}}$$
$$T = 4\sqrt{\frac{a}{g}} \int_0^{\pi/2} \frac{d\varphi}{\sqrt{1-k^2\sin^2\varphi}} = 4\sqrt{\frac{a}{g}} F\left(k, \frac{\pi}{2}\right)$$

$$F(k,\varphi) = \int_{0}^{\varphi} \frac{d\varphi}{\sqrt{1-k^{2} \sin^{2} \varphi}}$$
The elliptic integral
of the first kind

$$E(k,\varphi) = \int_{0}^{\varphi} \sqrt{1-k^{2} \sin^{2} \varphi} d\varphi$$
The elliptic integral
of the second kind

$$T = 4\sqrt{\frac{a}{g}}F\left(k,\frac{\pi}{2}\right)^{k=0} T = 4\sqrt{\frac{a}{g}} \times \frac{\pi}{2}$$

$$T = 2\pi\sqrt{\frac{a}{g}}$$
Problems pages 25, 26

C



The Brachistochrone BRACHIS = SHORT CHRONOS = TIME

What is the path - curve - producing the shortest possible time for a particle to descend from a given point to another Point.

The shortest distance between two points is a line, but the descent of a weighted particle is acted upon, in the very least, by gravity.



The correct answer is that a body takes less time to fall along the arc of a circumference than to fall along the "line" of a corresponding chord. The cycloid path allows the particle to move rapidly at first, while in steep descent, and thus build up sufficient speed to overcome the greater distance the particle must travel. Thus, the speed of the descending particle is accelerated by gravity.







We now complete our discussion and discover what curve the brachistochrone actully is by solving this equation

$$dx = \left(\frac{y}{c-y}\right)^{1/2} dy$$

$$dx = \left(\frac{y}{c-y}\right)^{1/2} dy \qquad \left(\frac{y}{c-y}\right)^{1/2} = \tan \varphi \quad *$$

so that $y = c \sin^2 \varphi \qquad dy = 2c \sin \varphi \cos \varphi d\varphi$

 $dx = \tan \varphi dy = 2c \sin^2 \varphi d\varphi = c(1 - \cos 2\varphi)d\varphi$

Integration now yield

 $x = \frac{c}{2} (2\varphi - \sin 2\varphi) + c_1 \quad \text{Our curve is to pass through the} \\ x = y = 0 \quad \text{when} \quad \varphi = 0 \rightarrow c_1 = 0 \\ x = \frac{c}{2} (2\varphi - \sin 2\varphi)$



A cycloid is the locus of a point on the circumference of a circle rotating along a fixed line . . .

the curve of shortest time is a cycloid

One arch of the cycloid comes from one rotation of the circle and so is described by $0 \le \theta \le 2\pi$.

Some properties of cycloid:

(1) A particle slides along the curve from point *A* to a lower point *B* not directly beneath *A*. Among all possible curves joining *A* to *B*, the particle will take the least time if the curve is an inverted arch of a cycloid.

If we take one-half of a cycloid and turn it upside-down we get the brachistochrone for bead.





(2) Notice that it is vertical at the start to get up lots of initial speed and then flattens out at the end. This path has another interesting property, namely if the ride could somehow be started from rest at point B or C, the ride would last the same length of time as if it started from point A.





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